Effect of Fractional Fourier Transformation on Time-Frequency Distributions Belonging to the Cohen Class

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Abstract— We consider the Cohen class of time-frequency distributions, which can be obtained from the Wigner distribution by convolving it with a kernel characterizing that distribution. We show that the time-frequency distribution of the fractional Fourier transform of a function is a rotated version of the distribution of the original function, if the kernel is rotationally symmetric. Thus, the fractional Fourier transform corresponds to rotation of a relatively large class of time-frequency representations (phase-space representations), confirming the important role this transform plays in the study of such representations.

THE fractional Fourier transform [1]–[3] has found many applications in quantum mechanics and optics [1], [2], [4]–[8], and signal processing [2], [3], [9]–[12], [23]. The fractional Fourier transform has been related to wavelet transforms [2], [13], neural networks [13], and is also related to various chirp-related operations [2], [14], [15]. It can be optically realized much like the usual Fourier transform [2], [4], [5], [7], and can be simulated with a fast digital algorithm [9], [16]. Other applications which are currently under study, or which have been suggested, include phase retrieval, signal detection, radar, tomography, and data compression.

The ath-order fractional Fourier transform $x_a(t)$ of the function x(t) may be defined for 0 < |a| < 2 as

$$x_a(t) \equiv \int_{-\infty}^{\infty} B_a(t, t') x(t') dt'$$
$$B_a(t, t') \equiv \frac{\exp\left(-i\pi \operatorname{sgn}\left(\sin\phi\right)/4 + i\phi/2\right)}{|\sin\phi|^{1/2}} \cdot \exp\left[i\pi(t^2 \cot\phi - 2tt' \csc\phi + t'^2 \cot\phi)\right] \quad (1)$$

where $\phi \equiv a\pi/2$. The kernel approaches $B_0(t, t') \equiv \delta(t-t')$ and $B_{\pm 2}(t, t') \equiv \delta(t+t')$ for a = 0 and $a = \pm 2$, respectively.

The Wigner distribution $W_x(t, f)$ of a signal x can be defined in terms of the time-domain representation of that signal x(t) as [19], [20]

$$W_x(t, f) = \int_{-\infty}^{\infty} x(t + t'/2) x^*(t - t'/2) e^{-2\pi i f t'} dt'.$$
 (2)

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Roughly speaking, $W_x(t, f)$ is a function which gives the distribution of signal energy over time and frequency. It is possible to show that the Wigner distribution of $x_a(t)$ is a rotated version of that of x(t), a result independently arrived at by different authors [2], [5], [6], [3]. Let R_{ϕ} denote the operator which rotates a 2-D function counterclockwise by an angle ϕ , that is

$$\mathbf{R}_{\phi}[z(t, f)] = z(t\cos\phi + f\sin\phi, -t\sin\phi + f\cos\phi). \quad (3)$$

Then, the above-mentioned result may be stated as

$$W_{x_a}(t, f) = \mathbf{R}_{-\phi}[W_x(t, f)].$$
 (4)

A similar relation holds for the ambiguity function [2], [3]. A corollary of (4) is the following [2], [17]:

$$\{\mathcal{R}_{\phi}[W_x(t, f)]\}(t_a) = |x_a(t_a)|^2 \tag{5}$$

where \mathcal{R}_{ϕ} is the Radon transform operator. \mathcal{R}_{ϕ} takes the integral projection of the 2-D function $W_x(t, f)$ onto an axis making angle $\phi = a\pi/2$ with the t axis. We refer to this axis as the t_a axis, or the *a*th fractional Fourier domain. The concept of fractional Fourier domains is developed in [2] and [18].

In [3], Almeida further showed that a relation similar to (4) holds for (a modified form of) the short-time Fourier transform and spectrogram. This leads one to inquire whether a similar relation is valid for a more general class of time-frequency distributions. Here, we consider time-frequency distributions $T_x(t, f)$, which are members of the Cohen class. These distributions can be derived from the Wigner distribution through the relation [21], [22], [19]

$$T_x(t, f) = \int_{t'} \int_{f'} \Psi_T(t - t', f - f') W_x(t', f') \, dt' \, df'.$$
(6)

 $\Psi_T(t, f)$ is a kernel uniquely corresponding to the distribution T_x . We show below that T_x will satisfy a relation similar to (4) if the kernel function is rotationally symmetric around the origin, that is, if $\Psi(t, f)$ is a function of $(t^2 + f^2)^{1/2}$ only.

(The same condition can also be stated in terms of the alternative kernel function $\Phi(\theta, \tau)$, which is also employed in the study of the Cohen class [20], [19]. Since $\Phi(\theta, \tau)$ and $\Psi(t, f)$ constitute a 2-D Fourier transform pair, rotational symmetry of either implies rotational symmetry of the other.)

Thus, fractional Fourier transformation corresponds to rotation of not only the Wigner distribution, ambiguity function, and spectrogram, but a much larger class of time-frequency

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representations (phase-space representations). This not only confirms the important role this transform plays in the study of such representations, but also supports the notion of referring to the axis making angle $\phi = a\pi/2$ with the t axis as the ath fractional Fourier domain [2], [18]. The many applications of the fractional Fourier transform have been predominantly posed in terms of the Wigner distribution of the signal. With this result, it will be possible to deal with other time-frequency distributions which might be more appropriate for particular applications.

Theorem: Let the signal x have a time-frequency representation $T_x(t, f)$, which is a member of Cohen's class with kernel function $\Psi_T(t, f)$. Then,

$$T_{\boldsymbol{x}_a}(t,f) = \mathbf{R}_{-\phi}[T_{\boldsymbol{x}}(t,f)] \tag{7}$$

for all x and $\phi = a\pi/2$, if $\Psi_T(t, f)$ is rotationally symmetric around the origin.

Proof: We only sketch the main features of the proof since the operations involved are elementary. First, apply $R_{-\phi}$, as defined in (3), to both sides of (6) to obtain

$$\boldsymbol{R}_{-\phi}[T_x(t,f)] = \int_u \int_v \Psi_T(t\cos\phi - f\sin\phi - u, t\sin\phi + f\cos\phi - v)W_x(u,v)\,du\,dv.$$
 (8)

Now, consider an instance of (6) for the function $x_a(t)$, rather than x(t). Use (4) and (3) to replace $W_{x_a}(t', f')$ by $W_x(t'\cos\phi - f'\sin\phi, t'\sin\phi + f'\cos\phi)$. Finally, make the change of variables $u = t'\cos\phi - f'\sin\phi$, $v = t'\sin\phi + f'\cos\phi$, to obtain

$$T_{x_a}(t, f) = \int_u \int_v \Psi_T(t - u\cos\phi - v\sin\phi, f + u\sin\phi - v\cos\phi) W_x(u, v) \, du \, dv.$$
(9)

Equation (7) will be true if the right-hand sides of the last two equations are equal for all x and for all ϕ (or a). This will be the case if

$$\Psi_T(t\cos\phi - f\sin\phi - u, t\sin\phi + f\cos\phi - v) = \Psi_T(t - u\cos\phi - v\sin\phi, f + u\sin\phi - v\cos\phi) \quad (10)$$

for all of the appearing variables. Now, it is not difficult to show, by the transformation $t' = t - u \cos \phi - v \sin \phi$, $f' = f + u \sin \phi - v \cos \phi$, that this condition is equivalent to $\Psi_T(t, f)$ being rotationally symmetric. This completes the proof.

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